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## DISSOCIATION ENERGIES, FRANCK-CONDON FACTORS AND R-CENTROIDS OF THE SULPHIDES OF METALS (LI AND NI )

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#### Abstract

The experimental potential energy curves for the different electronic states of astrophysically important $\mathrm{LiS}^{+}$and NiS molecules are constructed by using Rydberg-Klein-Rees method as modified by Vanderslice et al. The ground state dissociation energies are determined by curve fitting techniques using the five parameters Hulburt-Hirschfelder ( $\mathrm{H}-\mathrm{H}$ ) function. The estimated dissociation energies are $0.81 \pm 0.02$ and $3.56 \pm 0.08 \mathrm{eV}^{2}$ for $\mathrm{LiS}^{+}$, and NiS respectively. These values are in good agreement with the values reported in the earlier literature. The r-centroids and Franck-Condon factors (FC factors) for the bands of $3 \Pi-\mathrm{X}^{3} \Sigma^{-}$system of $\mathrm{LiS}^{+}$and $\mathrm{A}^{3} \Pi-\mathrm{X}^{3} \Sigma^{-}$system for NiS , molecules have been calculated. The Franck-Condon factors are evaluated by the approximated analytical methods of Jarmain and Fraser. The absence of the bands of these systems is explained.


Keywords: Diatomic molecules; Dissociation energies; Franck-Condon factors; r-centriods

## I. INTRODUCTION

There has been considerable interest in spectroscopic studies of transition metal containing species because of their importance in chemistry, astrophysics, surface science and ab initio calculations. These compounds are widely used as catalysts in various chemical processes.
Theoretical calculations on these molecules often fail to predict the correct spectroscopic properties and energy ordering of electronic states. The spectroscopic studies of these molecules, therefore, have become important to gain insight into chemical bonding in simple
metal systems. Ram et al. [1] have mentioned that NiS might also be present in S-type stars.
There are several theoretical studies on the spectroscopic properties of low-lying electronic states of NiS. An experimental observation of NiS has been reported by Zheng et al. [2] in the $495-555 \mathrm{~nm}$ regions. Ram et al. [1] have reported on the observation of NiS bands in the near IR region (2500-7500 $\mathrm{cm}^{-1}$ ). Molecular species containing sulphur atom are of astrophysical importance. Some of them have been definitely identified. Theoretical investigations of the $\mathrm{SH}^{+}$and $\mathrm{LiS}^{+}$cations have been pursued by Khadri et al. [3]. In view of their astrophysical / scientific importance, it is proposed to construct potential curves and evaluate dissociation energies. As no attempt has been made so far to evaluate dissociation energies, Franck-Condon factors (FC factors)
and r-centroids using spectroscopic data for $\mathrm{LiS}^{+}$and NiS molecules, the present authors have taken up this investigation.

## II. THEORY \& COMPUTATIONAL PROCEDURE

## Potential energy curves

Experimental observed vibrational levels are used to construct the potential energy curve. The molecular constants required for the present study have been taken from the literature [1-4] and are listed in Table I. Many authors employed RKRV method and constructed the potential energy curves [5-16]. The RKRV [1721] method is an improved form of RKR method, which makes use of Wentzel-KramersBrillouin (WKB) approximation from which we obtain reliable potential energy curves with the observed vibrational and rotational constants. The potential energy curves are constructed using the method of RKRV [17-21] utilizing the molecular constants (listed in Table I) and the results are presented in Table II. The details of the said method were reported in literature, as such only the results of the present work are given Table I.

## Dissociation energy

The RKRV turning points are inserted into the five parameter Hulburt-Hirschfelder's [22-24] function and the potential energies $\mathrm{U}(\mathrm{r})$ are calculated by varying the dissociation energy De values. An average percentage deviation is determined between the calculated $\mathrm{U}(\mathrm{r})$ and the experimental $G(v)$ values. An accurate estimation of the dissociation energy ( $\mathrm{D}_{\mathrm{o}}$ ) requires an empirical potential function, which provides the best reproduction of the experimental energy values. A critical evaluation of the importance of these functions was given by Steele et al. [24] and they have shown that the potential function of HulburtHirschfelder (H-H) [22-24] fits well with the RKRV curves of a large number of diatomic molecules. In the present investigation, it is
observed that the $\mathrm{H}-\mathrm{H}$ function [22, 23] functions fits best and reproduces the experimental energy values (Table III). Different $\mathrm{D}_{\mathrm{e}}$ values are used in the $\mathrm{H}-\mathrm{H}$ function and the $D_{e}$ value, for which the best fit of the energy values $U(r)$ is observed, which is taken as the dissociation energy ( $\mathrm{D}_{\mathrm{e}}$ ) of the molecule.

## r-Centroids and Franck-Condon factors

The Franck-Condon factors depend mainly on the shape and relative separation in the internuclear distance between the upper and lower potentials involved. These factors control the form of the wave function $\psi_{v^{\prime}}, \psi_{v^{\prime \prime}}$ and their overlap. In order to calculate Franck-Condon factors, one has to evaluate the integral $\psi_{v^{\prime}}$ $\psi_{v^{\prime}}$ dr. In the absence of exact expression for $\psi_{v}$, one may use the eigen functions derived from same approximate potential functions. Therefore, the methods for calculating FranckCondon factors depend on the choice of approximate eigen functions $\psi_{v}$ obtained by solving Schroedinger equations with approximate potential functions.
Fraser and Jarmain [25] introduced the method of average ' $\alpha$ ' Morse constant, which allows analytic integration of the integral

$$
\begin{equation*}
\mathrm{q}_{v^{\prime}, v^{\prime \prime}}=\int_{0}^{\infty} \psi_{v^{\prime}} \mathrm{r}^{\mathrm{n}} \psi_{v^{\prime \prime}} \mathrm{dr} \tag{1}
\end{equation*}
$$

where $n=0,1,2$. The above integral becomes Franck-Condon factor when $\mathrm{n}=0$. This method successfully replaces cumbersome numerical integration for many band systems. The details of the method are described elsewhere [ 9,10 , 15, 25, 26 ]. The method yields accurate FranckCondon factor especially when vibrational quantum numbers are low. The r-Centroids $\mathrm{r}_{\mathrm{v}^{\prime}, v^{\prime \prime}}$ represent the characteristic internuclear separation of a $v^{\prime} \rightarrow v^{\prime \prime}$ transition in a diatomic molecule band system and are defined by

$$
\begin{equation*}
\mathrm{r}_{\mathrm{v}^{\prime}, v^{\prime \prime}}=\frac{\int \psi_{v^{\prime}} r \psi_{v^{\prime}} d r}{\int \psi_{v^{\prime}} \psi_{v^{\prime \prime}} d r} \tag{2}
\end{equation*}
$$

By use of the quadratic equation method of Nicholls and Jarmain [26] and Fraser and Jarmain [25, 26], r-Centroids are estimated for different bands.

## III. RESULTS AND DISCUSSION

The inherent error in the $\mathrm{H}-\mathrm{H}$ function given by Steele et al [24] is $2 \%$. The error involved in the evaluation of $D_{e}$ is minimum average percentage deviation plus the inherent error in the potential function. The $D_{0}$ values are found to be $0.79 \pm 0.02$ and $3.52 \pm 0.08 \mathrm{eV} \mathrm{LiS}^{+}$and NiS respectively. Here $D_{0}=D_{e}-G(0)$. The turning points obtained from RKRV [17-21] method are inserted in $\mathrm{H}-\mathrm{H}$ [22-24,27] function and evaluated $\mathrm{U}(\mathrm{r})$ values. The relevant $\mathrm{U}(\mathrm{r})$ values for the selected $D_{e}$ values are given in Table III. The estimated $\mathrm{D}_{0}$ value for $\operatorname{LiS}^{+}(0.79 \pm 0.02 \mathrm{eV})$ is in good agreement with the value $(0.83 \mathrm{eV})$ recommended by Boldrev et al. [28]. The dissociation energy for the ground state of NiS $(3.52 \pm 0.08 \mathrm{eV})$ is in close agreement with the value 3.53 eV quoted by Huber and Herzberg [29] and also the value ( 3.58 eV ) listed by Lide [30].
The vibrational sum rule is satisfied for $v^{\prime}$ and $v^{\prime \prime}$ progressions in the case of ${ }^{1} \Pi-X^{3} \Sigma^{-}$transition of $\mathrm{LiS}^{+}$. $(1,1),(3,2)$ and $(5,0)$ bands are found to be small in magnitude, which indicates that the bands are not experimentally be observed. From the magnitudes of the F-C factors of the ${ }^{1} \Pi$ $\mathrm{X}^{3} \sum^{-}$system of $\mathrm{LiS}^{+}$, one may conclude that in the $v^{\prime}$ progression with $v^{\prime}=0$ with $v^{\prime \prime} \geq 5$ and $v$ 'progression with $v^{\prime}=0 v^{\prime \prime} \geq 5$ would not be observed. As seen from Table IV, for ${ }^{1} \Pi-X^{3} \Sigma^{-}$ the order of magnitude of the F-C factors for the bands $(1,3),(3,1)$ and also $(0,2),(2,0)$ are equal which suggest that bands are of nearly equal intensity.
Keen observation of the Table IV reveals that in the case of $2^{1} \Sigma-\mathrm{X}^{3} \Sigma^{-}$system of $\mathrm{LiS}^{+}$that the vibrational sum rule is not satisfied for $v^{\prime}$ and
$v^{\prime \prime}$ progressions as far as the bands are observed. Therefore, it is suggested that if more experimental work is done on the above system, it would definitely yields more new bands. Among all bands observed in this system (5, 0) band is found to be the strongest. The magnitude of the F-C factors in the case of A-X and ${ }^{1} \Pi$-X systems of NiS beyond $(0,3)$ and $(3,0)$ are very small in comparison with the other values and hence they may not be observed experimentally [25].
The F-C factor for $(0,0)$ band is strongest in A-X system of NiS among all the bands observed. $(0,1),(1,0)$ and also ( 1,1 ),(2,1) bands of F-C factor values in A-X system of NiS are appears to have equal intensity.
To understand the physical conditions of the emitter from the relative intensities of the bands of a molecule. One must have a theoretical knowledge of the corresponding r-Centroids and vibrational transition probabilities for the respective band heads in a band system. In the present study, the r-Centroids for the ${ }^{3} \Pi-X^{3} \Sigma^{-}$ and $2^{1} \Sigma^{-} \mathrm{X}^{3} \Sigma^{-}$systems of $\mathrm{LiS}^{+}$are calculated by the quadratic equation method of Nicholls and Jarmain [26] are presented in Table IV. It is of interest to note that $\overline{\mathrm{r}}_{0,0}$ is greater than $\left(\mathrm{r}_{\mathrm{e} 1}+\mathrm{r}_{\mathrm{e} 2}\right) / 2$ in the case of $2^{1} \sum-\mathrm{X}{ }^{3} \sum^{-}$which shows that the potential functions are not very anharmonic. The $\Delta \overline{\mathrm{r}}$ values for the band system $\mathrm{A}-\mathrm{X}$ of $\mathrm{LiS}^{+}$is greater than $0.01 \mathrm{~A}^{0}$, which indicates that the potentials are wide. The sequence difference $\Delta \overline{\mathrm{r}}$ $=\bar{r}_{v^{\prime}+1}, \quad v^{\prime \prime}+1-\bar{r}_{v^{\prime}, v}$ has been found to be constant in accordance with the observations of Nicholls and Jarmain [26]. Since $\mathrm{r}_{\mathrm{e} 1}>\mathrm{r}_{\mathrm{e} 2}$ in all cases of $\mathrm{LiS}^{+}$band systems. $r$-Centroids estimates are expected to increase with wavelength, which is the trend observed in a red-degraded band system.
The values of r-Centroids of NiS molecule are estimated and presented in the Table IV. The
value of $\quad \overline{\mathrm{r}}_{\mathrm{v}^{\prime} \mathrm{v}^{\prime \prime}}$ (r-Centroids) was found to increase with the increase of the $\lambda_{v}$, $v^{\prime \prime}$, value as has been observed by Nicholls and Jarmain [26]. The internuclear distance of the upper state is greater than the ground state ( $\mathrm{r}_{\mathrm{e} 1}>\mathrm{r}_{\mathrm{e} 2}$ ) in all cases, r-Centriods estimates are expected to increase with wave lengths which is the trend observed in red-degraded band systems of NiS molecule. The sequence difference $\Delta \overline{\mathrm{r}}=\overline{\mathrm{r}}_{\mathrm{v}^{\prime}+1, \mathrm{v}^{\prime \prime}+1}-\overline{\mathrm{r}}_{\mathrm{v}^{\prime}, v^{\prime \prime}}$ for this system is found to be a constant $\left(0.007 \mathrm{~A}^{0}\right)$. It is seen that $\Delta \overline{\mathrm{r}}$ is less than $0.01 \mathrm{~A}^{0}$ for this system suggesting that the potentials are not wide. In the case of $\mathrm{A}^{3} \Pi_{i}$ $-\mathrm{X}^{3} \Sigma^{-}$and $[5.8]^{1} \Pi-\mathrm{X}^{3} \Sigma^{-}$system of $\mathrm{NiS}, \mathrm{r}_{0,0}$ is found to be slightly greater than $\left(r_{e_{1}}+\mathrm{r}_{\mathrm{e} 2}\right) / 2$ indicating there by that the potentials are not very anharmonic $[9,10]$

## REFERENCES

1. R.S. Ram, S. Yu, I. Gordon and P.F. Bernath, J. Mol. Spectr. 20, 258 (2009).
2. X.F. Zheng, T.T. Wang, J.R. Guo, C.X. Chen andY. Chen, Chem. Phys. Lett. 137, 394 (2004).
3. F. Khadri, H.N.Dome, S.Lahman, Z. Benhakder and M. Hochalf, J. Mol. Spectr. 232, 237 (2006).
4. K.D. Setzer, F. Meinecke and E.H. Fink, J. Mol. Spectr. 56, 258 (2009).
5. R.R.Reddy and R. Viswanath, J. Astrophys. Astro. 153, 10 (1989).
6. R.R.Reddy, Y.N. Ahammed, K. RamaGopal and D. Baba Basha, J. Quant. Spectr. Radiat. Trans. 85, 105 (2004).
7. B. Narasimhamurthy and N. Rajamanickam, J. Astrophys. Astro. 53, 4 (1993).
8. R.R.Reddy and R. Viswanath, Astrophys. Space Sci., 113, 164 (1990).
9. R.R.Reddy, Y. N. Ahammed, D. Sasikala Devi, K. RamaGopal, P. Abdul Azeem and T.V.R.Rao, Astro. Phys Space Sci. 281, 729 (2002).
10. R.R.Reddy, K. Rama Gopal, Y. N.Ahammed, D. Baba Basha, K. Narasimhulu and L.S.S.Reddy, Ind. J. Pure Appl. Phys. 43, 237 (2005).
11. R.R.Reddy, Y. N. Ahammed, A.S. Rao and T.V.R. Rao, J. Quan. Spectr. Radiat. Trans. 1035, 54 (1995).
12. P. Sambasiva Rao, R. R.Reddy and D. M. Potukchi, J. Quant. Spectr. Radiat. Trans. 98, 81 (2006).
13. R. R. Reddy, Y.N. Ahammed, K. R. Gopal, P.A. Azeem and S. Anjaneyulu, Astrophys. Space Sci., 223, 262 (1999).
14. R. R. Reddy, Y.N. Ahammed, K. R. Gopal, P.A. Azeem and T.V.R.Rao, J. Quan. Spectr. Radiat. Trans. 501, 66 (2000a).
15. R. R. Reddy, Y.N. Ahammed, K. R. Gopal and D. Baba Basha, Astro. Phys. Space Sci. 286, 419 (2003).
16. J.B. Bharitya, S.H. Behere and M.L.P. Rao, J. Quan. Spectr. Radiat. Trans., 95, 43 (1990).
17. R. Z. Rydberg, Phys. 376, 73 (1931).
18. O.Z. Klein, Phys. 226, 6 (1932).
19. A.G. Rees, Proc. Phys. Soc., 938, 59 (1947).
20. J.T. Vanderslice, E.A. Masan and W.G. Lippincott, J. Mol. Spectr. 17, 3 (1959).
21. J.T. Vanderslice, E.A. Masan and W.G. Lippincott, J. Mol. Spectr. 83, 5 (1960).
22. H.M. Hulburt and J.O. Hirschfelder, J. Chem. Phys. 61, 9 (1941).
23. H.M. Hulburt and J.O. Hirschfelder, J. Chem. Phys. 1901, 35 (1961).
24. D. Steele, E.R. Linpincott and J.T. Vanderslice, Rev. Mod. Phys. 239, 34 (1962).
25. P.A. Fraser and W.R. Jarmain, Proc. Phys. Soc. 1153, 66A (1953).
26. R.W. Nicholls and W.R. Jarmain, Proc. Phys. Soc. 253, 69A (1956).
27. A.G. Gaydon, Editor, Dissociation energies and spectra of diatomic molecules, $3^{\text {rd }}$ edition, Proc. Phys. Soc., 253, 69A (1968).
28. A. I.Boldrev, J. Simon and P.V. R. Senleyer, J. Chem. Phys. 99 (1993).
29. K.P. Huber and G. Hertzburg, Constants of diatomic molecules, Newyork, Nostrand, Reinhold, (1979).
30. D.R. Lide, Editor, CRC Hand Book of Chemistry and Physics, $74^{\text {th }}$ edition, Tokyo, CRC press, (1993).

Table I: Molecular constants for different electronic states of $\mathrm{LiS}^{+}$and NiS

| Molecule | $\mu$ | State | $\begin{array}{ll} \hline \mathrm{T}_{\mathrm{e}}{ }^{1}{ }^{1} \text { (cm } \\ \hline \end{array}$ | $\left.\begin{array}{c} \omega_{\mathrm{e}}\left(\mathrm{~cm}^{-1}\right. \\ \\ \\ \\ \\ \hline \end{array}\right)$ | $\begin{gathered} \omega_{\mathrm{e}} \mathrm{X}_{\mathrm{e}} \\ \left(\mathrm{~cm}^{-1}\right) \end{gathered}$ | $\left.\begin{array}{\|cc\|} \hline \mathrm{B}_{\mathrm{e}}\left(\mathrm{ccm}^{-}\right. \\ & \\ & 1 \end{array}\right)$ | $\left.\begin{array}{lll} \hline \alpha_{e} & \left(\mathrm{~cm}^{-}\right. \\ & & { }^{1} \end{array}\right)$ | $\mathrm{re}_{\mathrm{e}}(\AA)$ | Reference |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{LiS}^{+}$ | 5.7059 | $\mathrm{X}^{3} \Sigma^{-}$ | 0 | 376.1 | 4.50 | 0.562 | 0.011 | 2.43 | Khadri al. ${ }^{8}$ | et |
|  |  | 3П | 5597.5 | 226.4 | 9.63 | 0.496 | 0.021 | 2.59 |  |  |
|  |  | $1 \Delta$ | 8904.5 | 385.5 | 4.62 | 0.566 | 0.008 | 2.42 |  |  |
|  |  | $1^{1} \Sigma^{+}$ | 14171.43 | 322.3 | 16.19 | 0.546 | 0.015 | 2.47 |  |  |
|  |  | $1 \Pi$ | 14518.26 | 210.9 | 4.17 | 0.480 | 0.023 | 2.63 |  |  |
|  |  | $2^{1} \Sigma$ | 27479.8 | 156.3 | 3.98 | 0.385 | 0.007 | 2.94 |  |  |
| NiS | 20.602 | $\mathrm{X}^{3} \Sigma^{-}$ | 0 | 512.68 | 3.61 | 0.212 | 0.00118 | 1.962 | $\mathrm{Ram}_{\text {al. }{ }^{1}}{ }^{1}$ | et |
|  |  | $\mathrm{A}^{3} \Pi_{\mathrm{i}}$ | 3864.56 | 466.69 | 1.1515 | 0.202 | 0.000728 | 2.012 |  |  |
|  |  | ${ }^{51.8]}{ }^{1} \pi$ | 5837.58 | 427.66 | 3.329 | 0.1956 | 0.001194 | 2.045 |  |  |

Table II: The turning points of the potential energy curves for the different electronic states of $\mathrm{LiS}^{+}$and NiS , molecules.

| LiS+ molecule |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{X}^{3} \Sigma^{-}$State, $\mathrm{T}_{\mathrm{e}}=0$ |  |  | $3 \Pi$ state, $\mathrm{T}_{\mathrm{e}}=5597.59 \mathrm{~cm}^{-1}$ |  |  | $1 \Delta$ state, $\mathrm{T}_{\mathrm{e}}=8904.5 \mathrm{~cm}^{-1}$ |  |  | $\begin{gathered} 1^{1} \Sigma^{+} \text {state, } \mathrm{T}_{\mathrm{e}}=14171.43 \\ \mathrm{~cm}^{-1} \end{gathered}$ |  |  | $2^{1} \sum$ State, $\mathrm{T}_{\mathrm{e}}=27479.8 \mathrm{~cm}^{-1}$ |  |  |
| $v$ | $\underset{\left(\mathrm{cm}^{-1}\right)}{\mathrm{U}}$ | $\begin{aligned} & \mathrm{r}_{\text {max }} \\ & (\AA \mathrm{A}) \\ & \hline \end{aligned}$ | $\begin{aligned} & \mathrm{r}_{\text {min }} \\ & (\AA \mathrm{A}) \end{aligned}$ | $\begin{gathered} \mathrm{U} \\ \left(\mathrm{~cm}^{-1}\right) \end{gathered}$ | $\begin{aligned} & r_{\text {max }} \\ & (\AA) \end{aligned}$ | $\mathrm{r}_{\text {min }}$ <br> (A) | $\underset{\left(\mathrm{cm}^{-1}\right)}{\mathrm{U}}$ | $\begin{aligned} & r_{\text {max }} \\ & (\AA \AA) \\ & \left(\begin{array}{l} \text { an } \end{array}\right. \end{aligned}$ | $\begin{aligned} & \mathrm{r}_{\text {min }} \\ & (\AA \mathrm{A}) \end{aligned}$ | $\underset{\left(\mathrm{cm}^{-1}\right)}{\mathrm{U}}$ | $\begin{aligned} & r_{\text {max }} \\ & (\AA) \\ & (\AA) \end{aligned}$ | $\begin{aligned} & \mathrm{r}_{\text {min }} \\ & (\AA \mathrm{A}) \end{aligned}$ | $\begin{gathered} \mathrm{U} \\ \left(\mathrm{~cm}^{-1}\right) \end{gathered}$ | $\mathrm{r}_{\text {max }}(\AA)$ | $\mathrm{r}_{\text {min }}(\mathrm{A})$ |
| 0 | 186.92 | 2.575 | 2.308 | 110.94 | 2.788 | 2.440 | 191.59 | 2.561 | 2.298 | 157.10 | 2.633 | 2.339 | 77.15 | 3.165 | 2.747 |
| 1 | 554.02 | 2.699 | 2.232 | 318.38 | 2.972 | 2.352 | 567.85 | 2.679 | 2.218 | 447.02 | 2.785 | 2.253 | 225.49 | 3.359 | 2.620 |
| 2 | 912.12 | 2.794 | 2.185 | 506.56 | 3.143 | 2.298 | 933.87 | 2.783 | 2.167 | 704.56 | 2.917 | 2.193 | 365.87 | 3.513 | 2.537 |
| 3 | 1261.22 | 2.879 | 2.150 | 675.48 | 3.303 | 2.257 | 1292.65 | 2.847 | 2.127 | 929.72 | 3.048 | 2.141 | 498.29 | 3.654 | 2.470 |
| 4 | 1601.32 | 2.957 | 2.122 | 825.14 | 3.469 | 2.221 | 1641.19 | 2.920 | 2.095 | 1122.50 | 3.188 | 2.090 | 622.75 | 3.790 | 2.410 |
| 5 | 1932.42 | 3.033 | 2.099 | 955.54 | 3.648 | 2.187 | 1980.49 | 2.989 | 2.669 | 1282.90 | 3.344 | 2.035 | 739.25 | 3.925 | 2.360 |
| 6 | 2254.52 | 3.106 | 2.079 | 1066.68 | 3.847 | 2.151 | 2310 | 3.056 | 2.042 | 1410.92 | 3.532 | 1.972 | 847.79 | 4.063 | 2.3119 |


| NiS molecule |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{X}^{3} \sum^{-}, \mathrm{T}_{\mathrm{e}}=0$ |  |  | $\mathrm{A}^{3} \Pi_{\mathrm{i}}, \mathrm{T}_{\mathrm{e}}=3864.56 \mathrm{~cm}^{-1}$ |  |  | [5.8] ${ }^{1} \Pi, \mathrm{~T}_{\mathrm{e}}=5837 \mathrm{~cm}^{-1}$ |  |  |
| $v$ | $\mathrm{U}^{\left.{ }^{1}\right)}{ }^{\left(\mathrm{cm}^{-}\right.}$ | $\mathrm{r}_{\text {max }}(\mathrm{A})$ | $\mathrm{r}_{\text {min }}(\mathrm{A})$ | $\begin{gathered} \mathrm{U} \\ \left(\mathrm{~cm}^{-1}\right) \end{gathered}$ | $\mathrm{r}_{\text {max }}(\mathrm{A})$ | $\mathrm{r}_{\text {min }}(\AA)$ | $\begin{gathered} \mathrm{U} \\ \left(\mathrm{~cm}^{-1}\right) \end{gathered}$ | $\mathrm{r}_{\text {max }}(\AA)$ | $\mathrm{r}_{\text {min }}(\mathrm{A})$ |
| 0 | 255.43 | 2.021 | 1.908 | 233.06 | 2.073 | 1.955 | 213.00 | 2.110 | 1.985 |
| 1 | 760.89 | 2.069 | 1.871 | 697.45 | 2.121 | 1.915 | 634.01 | 2.162 | 1.945 |
| 2 | 1259.13 | 2.104 | 1.847 | 1159.54 | 2.156 | 1.889 | 1048.36 | 2.200 | 1.919 |
| 3 | 1750.15 | 2.133 | 1.828 | 1619.33 | 2.184 | 1.869 | 1456.06 | 2.234 | 1.899 |
| 4 | 2233.95 | 2.160 | 1.812 | 2076.81 | 2.210 | 1.852 | 1857.09 | 2.264 | 1.881 |
| 5 | 27110.53 | 2.186 | 1.798 | 2532.00 | 2.23 | 1.837 | 2251.47 | 2.292 | 1.866 |
| 6 | 3179.89 | 2.209 | 1.786 | 2984.87 | 2.256 | 1.823 | 2639.19 | 2.318 | 1.852 |

Table III: Energy values using Hulbert - Hirschfelder (H-H) function for the $X^{\mathbf{3}} \sum^{\text {- }}$ electronic state of $\mathrm{LiS}^{+}$molecule

| $\left(\mathrm{A}^{\mathrm{o}}\right)$ | $\mathrm{U}\left(\mathrm{cm}^{-1}\right)$ | Calculated U(r)in cm |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | $\mathrm{D}_{\mathrm{e}}=6344 \mathrm{~cm}^{-1}$ | $\mathrm{D}_{\mathrm{e}}=6544 \mathrm{~cm}^{-1}$ | $\mathrm{D}_{\mathrm{e}}=6744 \mathrm{~cm}^{-1}$ |
| 2.575 | 186.9 | 182.01 | 187.74 | 193.48 |
| 2.698 | 554.0 | 539.57 | 556.58 | 573.59 |
| 2.794 | 912.1 | 888.24 | 916.24 | 944.24 |
| 2.878 | 1261.2 | 1227.61 | 1266.31 | 1305.01 |
| 2.957 | 1601.3 | 1557.25 | 1606.34 | 1655.44 |
| 3.032 | 1932.4 | 1876.75 | 1935.91 | 1995.08 |
| 3.105 | 2254.0 | 2185.66 | 2254.57 | 2323.47 |
| 2.308 | 186.9 | 181.95 | 187.22 | 193.48 |
| 2.231 | 554.0 | 538.59 | 555.57 | 572.55 |
| 2.184 | 912.1 | 884.90 | 912.80 | 940.69 |
| 2.149 | 1261.2 | 1220.27 | 1258.74 | 1297.21 |
| 2.121 | 1601.3 | 1544.20 | 1592.88 | 1641.56 |
| 2.098 | 1932.4 | 1856.29 | 1914.81 | 1973.33 |
| 2.078 | 2254.0 | 2156.25 | 2224.23 | 2292.21 |
| Average percentage deviation | 3.05 | 0.43 | 3.06 |  |

Energy values using Hulbert - Hirschfelder (H-H) function for the $X^{3} \Sigma^{-}$electronic state of NiS molecule

| $\left(\mathrm{A}^{\mathrm{o}}\right)$ | $\mathrm{U}\left(\mathrm{cm}^{-1}\right)$ | Calculated U(r) in cm |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | $\mathrm{D}_{\mathrm{e}}=25522 \mathrm{~cm}^{-1}$ | $\mathrm{D}_{\mathrm{e}}=28722 \mathrm{~cm}^{-1}$ | $\mathrm{D}_{\mathrm{e}}=28922 \mathrm{~cm}^{-1}$ |
| 2.021 | 255.43 | 253.33 | 255.11 | 256.88 |
| 2.069 | 760.89 | 754.21 | 759.50 | 764.79 |
| 2.104 | 1259.13 | 1248.59 | 1257.34 | 1266.10 |
| 2.133 | 1750.15 | 1732.16 | 1744.31 | 1756.45 |
| 2.160 | 2233.95 | 2209.12 | 2224.61 | 2240.10 |
| 2.186 | 2710.53 | 2678.14 | 2696.92 | 2715.70 |
| 2.209 | 3179.89 | 3139.27 | 3161.28 | 3183.29 |
| 1.908 | 255.43 | 255.11 | 255.20 | 256.98 |
| 1.871 | 760.89 | 759.50 | 760.82 | 766.12 |
| 1.847 | 1259.13 | 1257.34 | 1260.62 | 1269.40 |
| 1.828 | 1750.15 | 1744.31 | 1755.04 | 1767.26 |
| 1.812 | 2233.95 | 2224.61 | 2244.56 | 2260.20 |
| 1.798 | 2710.53 | 2696.92 | 2729.57 | 2748.57 |
| 1.786 | 3179.89 | 3161.28 | 3210.39 | 3232.75 |
| Average percentage deviation | 0.68 | 0.35 | 0.71 |  |

Table IV: Franck-Condon factors and r-centroids for the ${ }^{3} \Pi-X^{3} \sum^{-}$transition of $\mathrm{LiS}^{+}$molecule

| $\left(v^{\prime}, v^{\prime \prime}\right)$ | FC factor | r-centroids |
| :---: | :---: | :---: |
| $(0,0)$ | 0.495 | 2.528 |
| $(0,1)$ | 0.334 | 2.678 |
| $(0,2)$ | 0.132 | 2.813 |
| $(0,3)$ | 0.033 | 2.948 |
| $(0,4)$ | 0.005 | 3.088 |
| $(1,0)$ | 0.328 | 2.422 |
| $(1,1)$ | 0.029 | 2.592 |
| $(1,2)$ | 0.253 | 2.731 |
| $(1,3)$ | 0.247 | 2.862 |
| $(1,4)$ | 0.109 | 2.994 |
| $(1,5)$ | 0.027 | 3.132 |

Franck-Condon factors and r-centroids for the $2^{1} \sum-\mathrm{X}^{3} \sum^{-}$transition of $\mathrm{LiS}^{+}$molecule

| $\left(v^{\prime}, v^{\prime \prime}\right)$ | FC factor | r-centroids |
| :---: | :---: | :---: |
| $(0,0)$ | 0.004 | 2.696 |
| $(0,1)$ | 0.024 | 2.739 |
| $(0,2)$ | 0.070 | 2.780 |
| $(0,3)$ | 0.132 | 2.821 |
| $(0,4)$ | 0.183 | 2.861 |
| $(0,5)$ | 0.197 | 2.902 |
| $(1,0)$ | 0.016 | 2.679 |
| $(1,1)$ | 0.067 | 2.722 |
| $(1,2)$ | 0.123 | 2.763 |
| $(1,3)$ | 0.123 | 2.804 |
| $(1,4)$ | 0.060 | 2.844 |
| $(1,5)$ | 0.004 | 2.844 |
| $(2,0)$ | 0.035 | 2.662 |
| $(2,1)$ | 0.099 | 2.706 |
| $(2,2)$ | 0.103 | 2.747 |
| $(3,0)$ | 0.055 | 2.646 |
| $(3,1)$ | 0.102 | 2.690 |
| $(3,2)$ | 0.048 | 2.732 |
| $(4,0)$ | 0.012 | 2.631 |
| $(4,1)$ | 0.082 | 2.676 |
| $(5,0)$ | 0.819 | 2.616 |
| $(5,1)$ | 0.053 | 2.662 |
| $(5,2)$ | -- | 2.705 |
| $(5,3)$ | 0.068 | 2.745 |

Franck-Condon factors and r-centroids for the $\mathrm{A}^{3} \prod_{\mathrm{i}}-\mathrm{X}^{3} \sum^{-}$transition of NiS molecule

| $\left(v^{\prime}, v^{\prime \prime}\right)$ | FC factor | r-centroids |
| :---: | :---: | :---: |
| $(0,0)$ | 0.710 | 1.990 |
| $(0,1)$ | 0.249 | 2.063 |
| $(0,2)$ | 0.036 | 2.142 |
| $(0,3)$ | 0.003 | 2.230 |
| $(0,4)$ | -- | 2.328 |
| $(0,5)$ | 0.002 | 2.441 |
| $(1,0)$ | 0.230 | 1.929 |
| $(1,1)$ | 0.306 | 1.996 |


| $(1,2)$ | 0.357 | 2.069 |
| :--- | :--- | :--- |
| $(1,4)$ | 0.002 | 2.236 |
| $(1,5)$ | 0.095 | 2.334 |
| $(2,0)$ | 0.046 | 1.872 |
| $(2,1)$ | 0.310 | 1.935 |
| $(2,2)$ | 0.103 | 2.003 |
| $(3,0)$ | 0.007 | 1.819 |
| $(3,1)$ | 0.108 | 1.879 |
| $(3,2)$ | 0.579 | 1.942 |
| $(4,0)$ | 0.009 | 1.377 |
| $(4,1)$ | 0.032 | 1.825 |

Franck-Condon factors and r-centroids for the $[5.8]^{1} \Pi-X^{3} \sum^{-}$transition of NiS molecule

| $\left(v^{\prime}, v^{\prime \prime}\right)$ | FC factor | r-centroids |
| :---: | :---: | :---: |
| $(0,0)$ | 0.409 | 2.006 |
| $(0,1)$ | 0.389 | 2.053 |
| $(0,2)$ | 0.158 | 2.005 |
| $(0,3)$ | 0.036 | 2.161 |
| $(0,4)$ | 0.005 | 2.223 |
| $(0,5)$ | 0.001 | 2.292 |
| $(1,0)$ | 0.336 | 2.969 |
| $(1,1)$ | 0.004 | 2.013 |
| $(1,2)$ | 0.265 | 2.061 |
| $(1,3)$ | 0.268 | 2.113 |
| $(1,4)$ | 0.093 | 2.169 |
| $(1,5)$ | -- | 2.231 |
| $(2,0)$ | 0.163 | 1.935 |
| $(2,1)$ | 0.176 | 1.977 |
| $(2,2)$ | 0.065 | 2.021 |
| $(3,0)$ | 0.062 | 1.902 |
| $(3,1)$ | 0.206 | 1.943 |
| $(3,2)$ | 0.027 | 1.985 |
| $(3,3)$ | 0.013 | 2.029 |
| $(4,0)$ | 0.020 | 1.873 |
| $(4,1)$ | 0.135 | 1.911 |
| $(4,2)$ | 0.530 | 1.951 |

